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Hermite Expansion of the Riemann Zeta Function

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Abstract. Let $\zeta(s)$ be the Riemann zeta function, $s = \sigma + it$. For $0 < \sigma < 1$, we expand $\zeta(s)$ as the following series convergent in the space of slowly increasing distributions with variable *t*:

$$\zeta(\sigma+it) = \sum_{n=0}^{\infty} a_n(\sigma)\psi_n(t),$$

where

$$\psi_n(t) = (2^n n! \sqrt{\pi})^{-1/2} e^{\frac{-t^2}{2}} H_n(t),$$

 $H_n(t)$ is the Hermite polynomial, and

$$a_n(\sigma) = 2\pi(-1)^{n+1}\psi_n(i(1-\sigma)) + (-i)^n \sqrt{2\pi} \sum_{m=1}^{\infty} \frac{1}{m^{\sigma}}\psi_n(\ln m).$$

This paper is concerned with the convergence of the above series for $\sigma > 0$. In the deduction, it is crucial to regard the zeta function as Fourier transfomations of Schwartz' distributions.

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1 Results

Let $\zeta(s)$ be the famous Riemann Zeta function which is holomorphic for $s = \sigma + it \in C - \{1\}$. It is well known that if $0 < \sigma < 1$ then

$$\zeta(s) = s \int_0^\infty \frac{[x] - x}{x^{s+1}} dx$$

(see [1] or [2]). By the substitute of variables $x = e^y$, we get

$$\zeta(s) = s \int_{-\infty}^{\infty} ([e^y] - e^y) e^{-sy} dy.$$

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Set

$$f(y) = [e^y] - e^y.$$
(1.1)

Then for $0 < \sigma < 1$, $e^{-\sigma y} f(y)$ is a slowly increasing function, so can be regarded as an element of S', the dual space of the space S of rapidly decreasing functions on R. The Laplace transformation $\mathcal{L}(f)(s)$ of f is then defined on the trip $0 < \sigma < 1$ both in the ordinary and distributional sense, that is

$$\mathcal{L}(f)(s) = \zeta(s)/s.$$

Let $f' \in S'$ be the derivative of f in the distributional sense. then $e^{-\sigma y}f'(y) \in S'$ for $0 < \sigma < 1$, and Laplace transformation $\mathcal{L}(f')(s)$ is defined on the strip $0 < \sigma < 1$ such that

$$\mathcal{L}(f')(s) = \zeta(s)$$

([3] Chapter 8). So by the relation of Fourier and Laplace transformation of distribution (see also [3]), we see that $\zeta(\sigma+it)$ as function of t is the Fourier transformation of $e^{-\sigma y}f'(y)$ in the distributional sense, where the Fourier transformation is defined by $g(x) \rightarrow \int_{-\infty}^{\infty} g(x)e^{-ixy}dx$ for $g \in L^1(R)$.

Recall that the Hermite polynomials are defined as

$$H_n(x) = e^{x^2} \left(\frac{d}{dx}\right)^n e^{-x^2}, n = 0, 1, \dots$$

$$\psi_n(t) = \left(2^n n! \sqrt{\pi}\right)^{-1/2} e^{\frac{-t^2}{2}} H_n(t), n = 0, 1, \dots$$

form a complete normalized orthogonal system in $L^2(R)$.

Xiaqi Ding and his collaborators introduced and developed the theory of Hermite expansions of generalized functions [4]. The aim of this paper is to give the Hermite expansion of $\zeta(\sigma+it)$ as function of t for $0 < \sigma < 1$. For this, we give first the Hermite expansion of $e^{-\sigma y}f'(y) \in S'$ for fixed σ . Now

$$f'(y) = -e^y + \sum_{m=1}^{\infty} \delta(y - \ln m),$$

where δ is the Dirac δ -function. So

$$e^{-\sigma y} f'(y) = -e^{(1-\sigma)y} + \sum_{m=1}^{\infty} \frac{1}{m^{\sigma}} \delta(y - \ln m).$$
(1.2)

The following lemma gives the Hermite expansion of $-e^{(1-\sigma)y}$.

Lemma 1.1. For any complex number a,

$$\int_{-\infty}^{\infty} e^{ax} \psi_n(x) dx = (-i)^n \sqrt{2\pi} \psi_n(ia).$$

Proof. Since $e^{ax}\psi_n(x) \in L^1(R)$, its ordinary Fourier transformation exists. Let $c_n = (2^n n! \sqrt{\pi})^{\frac{1}{2}}$, then

$$\int_{-\infty}^{\infty} e^{ax} \psi_n(x) e^{-ixy} dx = \frac{(-1)^n}{c_n} \int_{-\infty}^{\infty} e^{\frac{x^2}{2} - i(y+ia)x} \left(\frac{d}{dx}\right)^n e^{-x^2} dx$$

$$= \frac{1}{c_n} \int_{-\infty}^{\infty} e^{-x^2} \left(\frac{d}{dx}\right)^n e^{\frac{x^2}{2} - i(y+ia)x} dx$$

$$= \frac{e^{\frac{1}{2}(y+ia)^2}}{c_n} \int_{-\infty}^{\infty} e^{-x^2} \left(\frac{d}{dx}\right)^n e^{\frac{1}{2}(x-i(y+ia))^2} dx$$

$$= \frac{i^n}{c_n} e^{\frac{1}{2}(y+ia)^2} \int_{-\infty}^{\infty} \left(\frac{d}{dy}\right)^n (e^{-x^2 + \frac{1}{2}(x-i(y+ia))^2} dx$$

$$= \frac{i^n}{c_n} e^{\frac{1}{2}(y+ia)^2} \left(\frac{d}{dy}\right)^n (e^{-\frac{1}{2}(y+ia)^2} \int_{-\infty}^{\infty} e^{-x^2 + \frac{1}{2}(x-i(y+ia))x} dx)$$

$$= \frac{i^n}{c_n} e^{\frac{1}{2}(y+ia)^2} \left(\frac{d}{dy}\right)^n (e^{-\frac{1}{2}(y+ia)^2} \sqrt{2\pi} e^{-\frac{1}{2}(y+ia)^2})$$

$$= \frac{(-i)^n \sqrt{2\pi}}{c_n} e^{-\frac{1}{2}(y+ia)^2} H_n(y+ia)$$

Taking y = 0 in the above formula, we get

$$\int_{-\infty}^{\infty} e^{ax} \psi_n(x) dx = (-i)^n \sqrt{2\pi} \psi_n(ia),$$

hence the lemma is proved.

Notice that although $e^{-\sigma y} f'(y) \in S'$, both the first and second terms in (1.2) are not in S'. In fact, they are only in \mathcal{D}' , the dual space of the compactly supported smooth functions. Especially, the series of the second term is only convergent in \mathcal{D}' . Since ψ_n is in S' but not in \mathcal{D}' , although it is easy to get

$$<\!\frac{1}{m^{\sigma}}\delta(y\!-\!\ln m),\psi_n(y)\!>=\!\frac{1}{m^{\sigma}}\psi_n(\ln m),$$

it is not obvious if the series

$$\sum_{m=1}^{\infty} \frac{1}{m^{\sigma}} \psi_n(\ln m) \tag{1.3}$$

converges. The following theorem addresses this issue.

Theorem 1.1. For $0 < \sigma < 1$ and $n = 0, 1, \dots$, the series (1.3) is convergent.

Proof. Set

$$\chi_1(y) = \begin{cases} 1, & y \in (-\infty, \ln \frac{3}{2}] \\ 0, & \text{otherwise} \end{cases}$$

and for $m = 2, 3, \cdots$,

$$\chi_m(y) = \begin{cases} 1, & y \in (\ln(m - \frac{1}{2}), \ln(m + \frac{1}{2})] \\ 0, & \text{otherwise.} \end{cases}$$

Denote by $g_m(y)$ the element of S' given by the integrable function $\chi_m(y)e^{-\sigma y}f(y)$. Then for any $\varphi \in S$, we have

$$< e^{-\sigma y} f(y), \varphi(y) > = \int_{-\infty}^{\infty} e^{-\sigma y} f(y) \varphi(y) dy$$

$$=\sum_{m=1}^{\infty}\int_{-\infty}^{\infty}\chi_m(y)e^{-\sigma y}f(y)\varphi(y)dy=\sum_{m=1}^{\infty}\langle g_m(y),\varphi(y)\rangle.$$

This means that

$$e^{-\sigma y}f(y) = \sum_{m=1}^{\infty} g_m(y)$$

is convergent in S'. Thus by a fundamental result in distributional theory, its derivative

$$(e^{-\sigma y}f(y))' = \sum_{m=1}^{\infty} g'_m(y)$$

is also convergent in \mathcal{S}' . Now

$$g'_{m}(y) = -\sigma g_{m}(y) + e^{-\sigma y} \chi'_{m}(y) f(y) + e^{-\sigma y} \chi_{m}(y) f'(y), \qquad (1.4)$$

$$e^{-\sigma y} \chi'_{1}(y) f(y) = -e^{-\sigma y} f(y) \delta\left(y - \ln\frac{3}{2}\right)$$

$$= -(\frac{3}{2})^{-\sigma} f\left(\ln\frac{3}{2}\right) \delta\left(y - \ln\frac{3}{2}\right), \qquad (1.4)$$

$$e^{-\sigma y} \chi'_{m}(y) f(y) = \left(m - \frac{1}{2}\right)^{-\sigma} f\left(\ln\frac{3}{2}\right) \delta\left(y - \ln(m - \frac{1}{2})\right)$$

$$- \left(m + \frac{1}{2}\right)^{-\sigma} f\left(\ln(m - \frac{1}{2})\right) \delta\left(y - \ln(m - \frac{1}{2})\right).$$

Since for any $\varphi \in S, \varphi(y) \rightarrow 0$ as $y \rightarrow \pm \infty$, and $|f(y)| \le 1$, we have

$$<\sum_{m=1}^{M} e^{-\sigma y} \chi'_{m}(y) f(y), \varphi(y) >$$
$$= -\left(M + \frac{1}{2}\right)^{-\sigma} f\left(\ln(M + \frac{1}{2})\right) \varphi\left(\ln(M + \frac{1}{2})\right) \to 0, \text{ as } M \to \infty.$$

Thus by (1.4), the series

$$\sum_{m=1}^{\infty} e^{-\sigma y} \chi_m(y) f'(y) = e^{-\sigma y} f'(y)$$
(1.5)

converges in \mathcal{S}' . Now

$$e^{-\sigma y}\chi_m(y)f'(y) = -e^{(1-\sigma)y}\chi_m(y) + \frac{1}{m^{\sigma}}\delta(y - \ln m)$$
(1.6)

and $e^{(1-\sigma)y}\psi_n(y)$ is integrable on *R*, so

$$\int_{-\infty}^{\infty} e^{(1-\sigma)y} \psi_n(y) dy = \sum_{m=1}^{\infty} \int_{-\infty}^{\infty} e^{-\sigma y} \chi_m(y) \psi_n(y) dy$$

is a convergent series. Thus by (1.6),

$$<\sum_{m=1}^{\infty}\frac{1}{m^{\sigma}}\delta(y-\ln m),\psi_n(y)>=\sum_{m=1}^{\infty}\frac{1}{m^{\sigma}}\psi_n(\ln m)$$

is a convergent series. The proof is complete.

By the general theory on Dirichlet series, we have **Corollary 1.1.** The Dirichlet series

$$\sum_{m=1}^{\infty} \frac{1}{m^s} \psi_n(\ln m)$$

is convergent for Re(s) > 0, and $n = 0, 1 \cdots$.

Corollary 1.2. For $0 < \sigma < 1$, the following Hermite expansion

$$e^{-\sigma y}f'(y) = \sum_{n=0}^{\infty} b_n(\sigma)\psi_n(y),$$

is convergent in \mathcal{S}' , where

$$b_n(\sigma) = -(-i)^n \sqrt{2\pi} \psi_n(i(1-\sigma)) + \sum_{m=1}^{\infty} \frac{1}{m^{\sigma}} \psi_n(\ln m).$$

Proof. Since $e^{-\sigma y} f'(y) \in S'$, it is proved in [4] that

$$e^{-\sigma y}f'(y) = \sum_{n=0}^{\infty} < e^{-\sigma y}f'(x), \psi_n(x) > \psi_n(y)$$

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holds in S'. By (1.5) and (1.6), we have

$$< e^{-\sigma x} f'(x), \psi_n(x) > = \sum_{m=1}^{\infty} < e^{-\sigma x} \chi_m(x) f'(x), \psi_n(x) >$$
$$= \sum_{m=1}^{\infty} \left(\int_{-\infty}^{\infty} e^{(1-\sigma)x} \chi_m(x) \psi_n(x) dx + \frac{1}{m^{\sigma}} \psi_n(\ln m) \right).$$

Thus Lemma 1.1 and Theorem 1.1 imply

$$b_n(\sigma) = < e^{-\sigma y} f'(x), \psi_n(x) >,$$

hence the corollary is proved.

Since

$$\int_{-\infty}^{\infty}\psi_n(x)e^{-ixy}dx=\sqrt{2\pi}(-i)^n\psi_n(y),$$

Corollary 1.2 implies that

Theorem 1.2. *For* $0 < \sigma < 1$,

$$\zeta(\sigma+it) = \sum_{n=0}^{\infty} a_n(\sigma)\psi_n(t)$$

in S', where

$$a_n(\sigma) = 2\pi(-1)^{n+1}\psi_n(i(1-\sigma)) + (-i)^n \sum_{m=1}^{\infty} \frac{\sqrt{2\pi}}{m^{\sigma}}\psi_n(\ln m).$$

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