

## VORTEX MOTION LAW OF AN EVOLUTIONARY GINZBURG-LANDAU EQUATION IN 2 DIMENSIONS\*

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**Abstract** We study the asymptotic behavior of solutions to an evolutionary Ginzburg-Landau equation. We also study the dynamical law of Ginzburg-Landau vortices of this equation under the Neuman boundary conditions. Away from the vortices, we use some measure theoretic arguments used by F.H.Lin in [1] to show the strong convergence of solutions. This is a continuation of our earlier work [2].

**Key Words** Ginzburg-Landau equations; vortex motion; asymptotic behavior;  $\epsilon$ -regularity.

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### 1. Introduction

We consider the following problem:

$$\frac{\partial u_\epsilon}{\partial t} = \Delta u_\epsilon + \frac{1}{\epsilon^2}(\beta^2(x) - |u_\epsilon|^2)u_\epsilon, \quad (x, t) \in \Omega \times R_+ \quad (1.1)$$

$$u_\epsilon(x, 0) = \beta u_\epsilon^0(x), \quad x \in \Omega \quad (1.2)$$

$$\frac{\partial u_\epsilon}{\partial \nu}(x, t) = 0, \quad x \in \partial\Omega, \quad t \geq 0 \quad (1.3)$$

where  $\Omega$  is a smooth bounded domain in  $R^2$ ,  $\nu$  the exterior unit normal vector along  $\partial\Omega$ .  $\beta(x) : \Omega \rightarrow R$  is a smooth function (say  $C^3$ ) with positive lower bound.  $u_\epsilon : \Omega \times R_+ \rightarrow R^2$ .

The initial datum  $\beta u_\epsilon^0(x)$  is smooth and satisfies (1.3). In addition, it also satisfies the following assumptions:

$$\|u_\epsilon^0(x)\|_{C(\bar{\Omega})} \leq K \quad (1.4)$$

$$\int_{\Omega} \rho^2(x) [|\nabla u_\epsilon^0(x)|^2 + \frac{1}{\epsilon^2} \beta^2(1 - |u_\epsilon^0|^2)^2] dx \leq K \quad (1.5)$$

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for a constant  $K$  and some  $m$  distinct points  $b_1, \dots, b_m$  in  $\Omega$ , where  $\rho(x) = \min\{|x - b_j| : j = 1, 2, \dots, m\}$ .

$$\begin{aligned} E(u_\varepsilon^0) &= \frac{1}{2} \int_{\Omega} [|\nabla u_\varepsilon^0|^2 + \frac{1}{2\varepsilon^2} \beta^2(x) (|u_\varepsilon^0(x)|^2 - 1)^2] dx \\ &\leq K[|\ln \varepsilon| + 1] \end{aligned} \quad (1.6)$$

The system (1.1)–(1.3) can be viewed as a simplified evolutionary Ginzburg-Landau equation in the theory superconductivity of inhomogeneity ([3]).

The aim of this article is to understand the dynamics of vortices, or zeros, of solutions  $u$  of (1.1)–(1.3). Its importance to the theory of superconductivity and applications is addressed in many earlier work ([3–7]).

Now we briefly describe some mathematical advances concerning this problem. In  $\beta = 1$ , the dynamical law for vortices was formally derived in [4,8]. The first rigorous mathematical proof of this dynamical law, which is of form  $\frac{\partial}{\partial t} a(t) = -\nabla w(a(t))$ , was given by F.H.Lin in [5,9]. See also [10, Lecture 3]. In [5,9], one allows the vortices of degree  $\pm 1$  and assumes that they have the same sign. For the vortices of degree  $\pm 1$  (possibly of different signs), the same dynamical law was shown later in [11]. We refer to [1] for vortex dynamics under the Neumann boundary conditions or pinning conditions. In the 3-dimensional case,  $\beta = 1$ , a similar dynamical law was also established in [1] for nearly parallel filaments. The short-time dynamical law for codimension 2 interfaces in higher dimensions was shown in [1]. When  $\beta \neq 1$ , in the 2-dimensional case, the dynamical law was established in [12] under the first boundary condition. But, here one proves only that  $u_\varepsilon$  converges weakly to the limit function in  $H_{loc}^1$  as  $\varepsilon \rightarrow 0^+$ .

The main goal of the present paper is to examine the vortex dynamics without topological constraints, and proves that  $u_\varepsilon$  converges strongly to the limit function in  $H_{loc}^1$  as  $\varepsilon \rightarrow 0^+$ .

To understand the behavior of  $u$  of (1.1)–(1.3) as  $t \rightarrow \infty$ , one has to look at steady state solutions  $u_\varepsilon$ , that is, the minimizer of the energy functional

$$E_\varepsilon(u) = \frac{1}{2} \int_{\Omega} \left[ |\nabla u|^2 + \frac{1}{2\varepsilon^2} (\beta^2 - |u|^2)^2 \right] dx$$

A complete characterization of asymptotic behavior (as  $\varepsilon \rightarrow 0^+$ ) of vortices of  $u_\varepsilon$  is given in the recent work [2].

Now we claim our main theorem.

**Theorem 1.1** *Assume that  $\beta \in C^3(\bar{\Omega})$  and  $\beta_0 = \min_{\bar{\Omega}} \beta(x) > 0$ . Under the assumptions (1.4)–(1.6), one has, for any  $0 \leq t \leq T$ , that*

$$u_\varepsilon(x, t) \rightarrow u_*(x, t) \quad (1.7)$$

*strongly in  $H_{loc}^1(\bar{\Omega} \times [0, T] \setminus \{(a_j(t), t) : t \in [0, T], j = 1, 2, \dots, m\})$ . Here the convergence is understood in the sense that for any sequence of  $\varepsilon'$  going to zero, there is a subsequence*